### 2. 2. Fin Performance

- ➤ The fin itself represents a conduction resistance to heat transfer from the original surface.
- > Accordingly, not always the extended surface enhances the heat transfer.
- The fin effectiveness  $(\varepsilon_f)$  is needed to evaluate the heat transfer through the extended surface.
- $\succ \varepsilon_f$  is defined as the ratio of the fin heat transfer rate to the heat transfer rate that would exist without the fin.



$$\varepsilon_f = \frac{q_f}{hA_{c,b}\theta_b}$$

where  $A_{c,b}$  is the fin cross-sectional area at the base.

### 2. 2. Fin Performance

- > The value of  $\varepsilon_f$  should be as large as possible, and in general, the use of fins may rarely be justified unless  $\varepsilon_f > 2$ .
- Based on the boundary conditions used for the investigated system, the effectiveness for a fin of uniform cross section may be calculated.
- ➤ assuming the convection coefficient of the finned surface to be equivalent to that of the base surface, it follows that, for the infinite fin approximation (Case D), the result is:  $(\mu p)^{1/2}$

$$\varepsilon_f = \left(\frac{kP}{hA_c}\right)^{1/2}$$

| TABLE 3.4 | Temperature | distribution a | and heat | loss for fi | ins of 1 | uniform | cross section |
|-----------|-------------|----------------|----------|-------------|----------|---------|---------------|
|-----------|-------------|----------------|----------|-------------|----------|---------|---------------|

| Case                                         | Tip Condition<br>(x = L)                                                                                                 | Temperature Distribution $\theta/\theta_j$ | Fin Heat<br>Transfer Rate $q_f$ |   |        |
|----------------------------------------------|--------------------------------------------------------------------------------------------------------------------------|--------------------------------------------|---------------------------------|---|--------|
| D                                            | Infinite fin $(L \rightarrow \infty)$ :<br>$\theta(L) = 0$                                                               | $e^{-mx}$                                  | (3.84)                          | М | (3.85) |
| $\theta \equiv T - T$ $\theta_b = \theta(0)$ | $ \begin{aligned} T_{\infty} & m^2 \equiv hP/kA_c \\ = T_b - T_{\infty} & M \equiv \sqrt{hPkA_c}\theta_b \end{aligned} $ |                                            |                                 |   |        |

### 2. 2. Fin Performance

- According to this equation, one can conclude that the following are needed achieving a fin-enhanced heat transfer:
- 1. High thermal conductivity (*Cu or Al*).
- 2. High ratio of the perimeter to the cross-sectional area: thin fins

TABLE 1.1

3. Fluids with low convection coefficient (gases).

 $k_{Cu} = 398 \text{ W/m.K}$  $k_{Al} = 180 \text{ W/m.K}$  $k_{Fe} = 14 \text{ W/m.K}$ 

| ansfer coefficient    | $\varepsilon_f = \left(\frac{kP}{L}\right)^1$              |
|-----------------------|------------------------------------------------------------|
| $h \ (W/m^2 \cdot K)$ | L = 2.65/L                                                 |
|                       |                                                            |
| 2–25                  |                                                            |
| 50-1000               |                                                            |
|                       |                                                            |
| 25-250                |                                                            |
| 100-20,000            |                                                            |
|                       | $\frac{h}{(W/m^2 \cdot K)}$ 2–25 50–1000 25–250 100–20,000 |

Typical values of the

### 2. 2. Fin Performance

- ▶ Fin performance may also be quantified in terms of a thermal resistance.
- > The *fin resistance* may be defined as:

$$R_{t,f} = \frac{\theta_b}{q_f}$$

> For the bas surface without fin, then. convection resistance will be considered:

$$R_{t,b} = \frac{1}{hA_{c,b}}$$

➤ Using those equations, the effectiveness can be then determined:

$$\varepsilon_f = \frac{R_{t,b}}{R_{t,f}}$$

To increase  $\varepsilon_f$  it is necessary to reduce *the conduction/convection resistance* of the fin.

$$\varepsilon_f = \left(\frac{kP}{hA_c}\right)^{1/2}$$

### 2. 2. Fin Performance

> Another measure of fin thermal performance is provided by the fin efficiency ( $\eta_f$ ).

$$\eta_f \equiv \frac{q_f}{q_{\text{max}}} = \frac{q_f}{hA_f\theta_b}$$

➢ For a straight fin of uniform cross section and an adiabatic tip:

$$\eta_f = \frac{M \tanh mL}{hPL\theta_b} = \frac{\tanh mL}{mL}$$

This result shows that η<sub>f</sub> approaches its maximum and minimum values of 1 and 0, respectively, as L approaches 0 and infinity.

### 2. 2. Fin Performance

#### Example-1

A stainless-steel pin fin that has k = 16 W/m.K, L = 10 cm, d = 1 cm and that is exposed to a boiling-water convection situation with h = 5000 W/m<sup>2</sup>.K. Based on these conditions, show whether this fin can enhance the heat transfer.



Thus, this rather large pin produces an increase of only 13 percent in the heat transfer.

- 4.1 Alternative Approaches
- 4.2 The Method of Separation of Variables
- 4.3 The Conduction Shape Factor and the Dimensionless Conduction Heat Rate
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References

Problems



- >Up to this point, conduction problems in which the temperature gradient is significant for only one coordinate direction were considered.
- > For more practical applications, multidimensional effects should be taken into account.



#### **1. Alternative Approaches**

- > Considering a prismatic solid in which there is two-dimensional heat conduction.
- > Two parallel surfaces are insulated at the same temperature, and the other surface were maintained at different temperatures (temperature gradient).
- > Considering, a steady state flow without energy generation.



FIGURE 4.1 Two-dimensional conduction.

#### 1. Alternative Approaches

- > The last equation can be solved using three various approaches giving the final equation to describe the temperature distribution along the two directions (x and y).
- □ Analytical: like variables separation and this can lead to an exact solution of the previous partial differential equation.
- Graphical: this method can provide only approximate results in a short time.
- Numerical (finite-difference, finite-element, or boundary-element): by this method, accurate results for complex, two- or three-dimensional geometries involving a wide variety of boundary conditions, can be reached.

#### 2. The Method of Separation of Variables

In order to obtain a simple solution; the transformation  $(\theta)$  was introduced.

$$\theta \equiv \frac{T - T_1}{T_2 - T_1}$$

$$\frac{\partial^2 \theta}{\partial x^2} + \frac{\partial^2 \theta}{\partial y^2} = 0$$

$$\partial^{2}T/\partial z^{2} \approx 0$$

$$\int_{W}^{y} T_{2}, \theta = 1$$

$$\int_{W}^{W} T_{2}, \theta = 1$$

$$T_{1}, \theta = 0$$

$$T_{1}, \theta = 0$$

$$\int_{U}^{U} T_{1}, \theta = 0$$

**FIGURE 4.2** Two-dimensional conduction in a thin rectangular plate or a long rectangular rod.

#### 2. The Method of Separation of Variables

Since the equation is second order in both x and y, two boundary conditions are needed for each of the coordinates.

$$\theta(0, y) = 0$$
 and  $\theta(x, 0) = 0$   
 $\theta(L, y) = 0$  and  $\theta(x, W) = 1$ 

Considering the solution is the product of two functions (one for and another for Y).

$$\theta(x, y) = X(x) \cdot Y(y)$$
$$-\frac{1}{X} \frac{d^2 X}{dx^2} = \frac{1}{Y} \frac{d^2 Y}{dy^2}$$



**FIGURE 4.2** Two-dimensional conduction in a thin rectangular plate or a long rectangular rod.

#### 2. The Method of Separation of Variables

> Using the separation constant as  $\lambda^2$ :

$$-\frac{1}{X}\frac{d^2X}{dx^2} = \frac{1}{Y}\frac{d^2Y}{dy^2} = \lambda^2$$

$$\frac{d^2 X}{dx^2} + \lambda^2 X = 0$$
$$\frac{d^2 Y}{dy^2} - \lambda^2 Y = 0$$

The last two equation are ordinary differential equations (one variable in each).



**FIGURE 4.2** Two-dimensional conduction in a thin rectangular plate or a long rectangular rod.

#### 2. The Method of Separation of Variables

> The solution of the pervious equation is:

 $X = C_1 \cos \lambda x + C_2 \sin \lambda x$  $Y = C_3 e^{-\lambda y} + C_4 e^{+\lambda y}$ 

> Then the general solution is:

$$\theta = (C_1 \cos \lambda x + C_2 \sin \lambda x)(C_3 e^{-\lambda y} + C_4 e^{\lambda y})$$

C1, C2, C3 and C4 are integration constant which can be determined using the boundary conditions (4 conditions) selected in the system.



**FIGURE 4.2** Two-dimensional conduction in a thin rectangular plate or a long rectangular rod.

2. The Method of Separation of Variables

 $\theta = (C_1 \cos \lambda x + C_2 \sin \lambda x)(C_3 e^{-\lambda y} + C_4 e^{\lambda y})$ <u>1. x = 0 or y = 0:</u>  $\theta(x, 0) = 0$   $C_2 \sin \lambda x (C_3 + C_4) = 0$  $C_2 = 0$ , this will result in: The case, then, is that  $C_3 + C_4$  $\theta(x, y) = 0$ = 0 and it does not satisfy the *boundary condition;*  $\theta(x, W) = 1$ 

- 2. The Method of Separation of Variables
- 2. Using the condition in which:  $\theta(L, y) = 0$

 $C_{4}$  cannot be zero.

$$C_{2}C_{4}\sin\lambda L(e^{\lambda y} - e^{-\lambda y}) = 0 \quad \text{in } \lambda L = 0$$

$$\lambda = \frac{n\pi}{n} \quad n = 1, 2, 3, \dots$$

$$n = \frac{n\pi}{L} \qquad n = 1, 2, 3, \dots$$

n cannot be zero.

2. The Method of Separation of Variables3. Using the condition in which:  $\theta(x, W) = 1$ 

$$1 = \sum_{n=1}^{\infty} C_n \sin \frac{n\pi x}{L} \sinh \frac{n\pi W}{L}$$

By applying a standard method,  $C_n$  can be determined finally using this equation:

$$C_n = \frac{2[(-1)^{n+1} + 1]}{n\pi \sinh(n\pi W/L)} \qquad n = 1, 2, 3, \dots$$

$$\theta(x, y) = \frac{2}{\pi} \sum_{n=1}^{\infty} \frac{(-1)^{n+1} + 1}{n} \sin \frac{n\pi x}{L} \frac{\sinh(n\pi y/L)}{\sinh(n\pi W/L)}$$

2. The Method of Separation of Variables

$$\theta(x, y) = \frac{2}{\pi} \sum_{n=1}^{\infty} \frac{(-1)^{n+1} + 1}{n} \sin \frac{n\pi x}{L} \frac{\sinh(n\pi y/L)}{\sinh(n\pi W/L)}$$





# 3. The Conduction Shape Factor and the Dimensionless Conduction Heat Rate

- In general, finding analytical solutions to the two- or three-dimensional heat equation is time-consuming and, in many cases, not possible.
- > Therefore, a different approach is often taken.
- This can be done using a shape factor S or a steady-state dimensionless conduction heat rate, q\*<sub>ss</sub>:

$$q = Sk\Delta T_{1-2}$$

> Where  $T_{1-2}$  is the temperature difference between boundaries.

$$R_{\rm t,cond(2D)} = \frac{1}{Sk}$$



# 3. The Conduction Shape Factor and the Dimensionless Conduction Heat Rate

| System                                                                                              | Schematic           | Restrictions                   | Shape Factor                                                 |
|-----------------------------------------------------------------------------------------------------|---------------------|--------------------------------|--------------------------------------------------------------|
| Case 1<br>Isothermal sphere buried in a semi-<br>infinite medium                                    |                     | z > D/2                        | $\frac{2\pi D}{1 - D/4z}$                                    |
| <b>Case 2</b><br>Horizontal isothermal cylinder of length <i>L</i> buried in a semi-infinite medium | $T_2$               | $L \gg D$ $L \gg D$ $z > 3D/2$ | $\frac{2\pi L}{\cosh^{-1}(2z/D)}$ $\frac{2\pi L}{\ln(4z/D)}$ |
| Case 3<br>Vertical cylinder in a semi-infinite<br>medium                                            | $T_2$<br>$T_1$<br>D | $L \gg D$                      | $\frac{2\pi L}{\ln\left(4L/D\right)}$                        |

## 3. The Conduction Shape Factor and the Dimensionless Conduction Heat Rate

| Physical system                                                                                       | Schematic                                                                                           | Shape factor                                                                          | Restrictions                                   |
|-------------------------------------------------------------------------------------------------------|-----------------------------------------------------------------------------------------------------|---------------------------------------------------------------------------------------|------------------------------------------------|
| Thin rectangular plate of<br>length L, buried in semi-infinite<br>medium having isothermal<br>surface | Isothermal $D \longrightarrow U$                                                                    | $\frac{\pi W}{\ln(4 W/L)}$ $\frac{2\pi W}{\ln(4 W/L)}$ $\frac{2\pi W}{\ln(2\pi D/L)}$ | D = 0<br>W > L<br>$D \gg W$<br>W > L<br>D > 2W |
| Parallel disks buried in<br>infinite medium                                                           | $ \begin{array}{ccc} t_1 & t_2 \\  & O \\  & O \\  & & & \\  & & & \\  & & & \\  & & & \\  & & & &$ | $\frac{4\pi r}{\left[\frac{\pi}{2} - \tan^{-1}(r/D)\right]}$                          | D > 5r<br>$\tan^{-1}(r/D)$<br>in radians       |
| Eccentric cylinders of length L                                                                       |                                                                                                     | $\frac{2\pi L}{\cosh^{-1}\left(\frac{r_1^2 + r_2^2 - D^2}{2r_1 r_2}\right)}$          | $L \gg r_2$                                    |
| Cylinder centered in a square<br>of length <i>L</i>                                                   |                                                                                                     | $\frac{2\pi L}{\ln(0.54 \text{ W/r})}$                                                | $L \gg W$                                      |

# 3. The Conduction Shape Factor and the Dimensionless Conduction Heat Rate

#### **Example-1**:

The possible existence of an optimum insulation thickness for radial systems is suggested by the presence of competing effects associated with an increase in this thickness. In particular, although the conduction resistance increases with the addition of insulation, the convection resistance decreases due to increasing outer surface area. Hence there may exist an insulation thickness that minimizes heat loss by maximizing the total resistance to heat transfer. Resolve this issue by considering the following system.



## 3. The Conduction Shape Factor and the Dimensionless Conduction Heat Rate

**Example-1**:



3. The Conduction Shape Factor and the Dimensionless Conduction Heat Rate

### **Example-1**:

- 1. Steady-state conditions.
- 2. One-dimensional heat transfer in the radial (cylindrical) direction.
- **3.** Negligible tube wall thermal resistance.
- 4. Constant properties for insulation.
- **5.** Negligible radiation exchange between insulation outer surface and surroundings.



3. The Conduction Shape Factor and the Dimensionless Conduction Heat Rate

**Example-1**:



For optimum insulation thickness

$$\frac{dR'_{\text{tot}}}{dr} = 0 \quad \implies \quad \frac{1}{2\pi kr} - \frac{1}{2\pi r^2 h} = 0 \qquad r = \frac{k}{h}$$

# 3. The Conduction Shape Factor and the Dimensionless Conduction Heat Rate

#### **Example-1**:

A metallic electrical wire of diameter d = 5 mm is to be coated with insulation of thermal conductivity k=0.35 W/m.K. It is expected that, for the typical installation, the coated wire will be exposed to conditions for which the total coefficient associated with convection and radiation is h = 15 W/m<sup>2</sup>.K. To minimize the temperature rise of the wire due to ohmic heating, the insulation thickness is specified so that the *critical insulation radius* is achieved (as shown in the previous example). During the wire coating process, however, the insulation thickness sometimes varies around the wire. Determine the change in the thermal resistance of the insulation due to the variation that is 50% of the critical insulation thickness.

3. The Conduction Shape Factor and the Dimensionless Conduction Heat Rate

### **Example-1**:

- 1. Steady-state conditions.
- 2. Two-dimensional conduction.
- 3. Constant properties.
- 4. Both the exterior and interior surfaces of the coating are at uniform temperatures.



# 3. The Conduction Shape Factor and the Dimensionless Conduction Heat Rate

**Example-1**:

$$r_{\rm cr} = \frac{k}{h} = \frac{0.35 \text{ W/m} \cdot \text{K}}{15 \text{ W/m}^2 \cdot \text{K}} = 0.023 \text{ m} = 23 \text{ mm}$$

the critical insulation thickness is

$$t_{\rm cr} = r_{\rm cr} - d/2 = 0.023 \text{ m} - \frac{0.005 \text{ m}}{2} = 0.021 \text{ m} = 21 \text{ mm}$$

The thermal resistance of the coating associated with the concentric wire

$$R'_{t,\text{cond}} = \frac{\ln[r_{\text{cr}}/(d/2)]}{2\pi k} = \frac{\ln[0.023 \text{ m}/(0.005 \text{ m}/2)]}{2\pi (0.35 \text{ W/m} \cdot \text{K})} = 1.0 \text{ m} \cdot \text{K/W}$$





# 3. The Conduction Shape Factor and the Dimensionless Conduction Heat Rate

#### **Example-1**:

For the eccentric wire, the thermal resistance of the insulation may be evaluated using



# 3. The Conduction Shape Factor and the Dimensionless Conduction Heat Rate

### **Example-2:**

Horizontal pipe 15 cm in diameter and 4 m long is buried in the earth at a depth of 20 cm. The pipe-wall temperature is 75 °C, and the earth surface temperature is 5 °C. Assuming that the thermal conductivity of the earth is 0.8 W/m. °C, calculate the heat lost by the pipe.



Since D < 3r, We may calculate the shape factor for this situation using the equation given in this table.

## 3. The Conduction Shape Factor and the Dimensionless Conduction Heat Rate

**Example-2:** 

$$S = \frac{2\pi L}{\cosh^{-1}(D/r)} = \frac{2\pi(4)}{\cosh^{-1}(20/7.5)} = 15.35 \text{ m}$$

The heat flow is calculated from

$$q = kS\Delta T = (0.8)(15.35)(75 - 5) = 859.6$$
 W